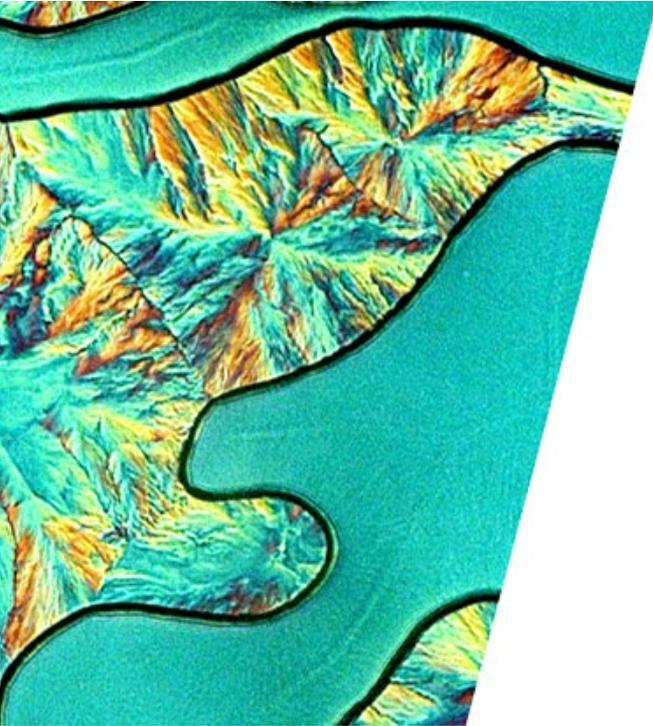


# The OSA Nonlinear Optics Technical Group Welcomes You!



High-Order Dispersion Solitons  
and Topological Photonics in  
Silicon

Andrea Blanco-Redondo

# Technical Group Leadership 2021



Chair

**Amol Choudhary**

Indian Institute of Technology, India



Vice Chair

**Ajanta Barh**

ETH Zürich, CH



Communications Officer

**Lin Xu**

Univ. of Southampton, UK



Webinar Officer

**Alexander Solntsev**

University of Technology Sydney, AU



Events Officer

**Donnie Keathley**

RLE, MIT, USA



Nonlinear  
Optics  
Technical Group

# Technical Group at a glance

- Focus
  - “Physics of nonlinear optical materials, processes, devices, & applications”
  - **3800** members (**largest** in OIS, 3<sup>rd</sup> largest in OSA)
- Mission
  - To benefit YOU
  - webinars, e-Presence, publications, technical events, business events, outreach
  - Interested in presenting your research? Have ideas for TG events? Contact us at
- Email: [TGNonlinearOptics@osa.org](mailto:TGNonlinearOptics@osa.org)
- Contact
  - [www.osa.org](http://www.osa.org)
  - Facebook: [www.facebook.com/osanonlineoptics](http://www.facebook.com/osanonlineoptics)
  - LinkedIn: [www.linkedin.com/groups/8302249](http://www.linkedin.com/groups/8302249)

# Today's webinar



## High-Order Dispersion Solitons and Topological Photonics in Silicon

Short bio:

- 2019 – now: Head of the Silicon Photonics department at Nokia Bell Labs
- 2015 – 2019: Professor Harry Messel Research Fellow and Senior Lecturer at the School of Physics of the University of Sydney, Australia
- 2007 – 2015: photonics researcher and a project manager with the Aerospace and Telecom departments of Tecnalia, Spain
- 2014: PhD in photonics at the University of the Basque Country, Bilbao, Spain
- OSA Director at Large, Member of the OSA Finance Council, Associate Editor at OSA Continuum

# High-order dispersion solitons and topological photonics in silicon

**Andrea Blanco-Redondo**

Nokia Bell Labs, New Providence, NJ, USA

[andrea.blanco-redondo@nokia-bell-labs.com](mailto:andrea.blanco-redondo@nokia-bell-labs.com)

# Collaborators

## Bell Labs (NJ, US):

Rene Essiambre, Nick Fontaine, Mark Earnshaw, Brian Stern, KW Kim, Mikael Mazur ...

## University of Sydney (Sydney, Australia):

Martijn de Sterke, Antoine Runge, Bryn Bell, Wei-Wei Zhang, Steven Bartlett, Ben Eggleton, Chad Husko, Alessio Stefani

*Students:* Michelle Wang, Cooper Doyle, Sam Lo, Kevin Tam, Joshua Lourdesamy, Imanol Andonegui, Ezgi Sahin

## Macquarie University (Sydney, Australia):

Darren Hudson (now CACI), Matthew Collins (now Xanadu)

## Technion (Haifa, Israel):

Moti Segev, Dikla Oren, Gal Harari, Kobi Lumer, Mik Rechtsman (now Penn State)

## University of York (York, UK):

Thomas Krauss, Juntao Li

## Soliton work



Prof. M. de Sterke  
*Usyd*

Dr. Antoine Runge  
*Usyd*

Dr. Darren Hudson  
*CACI - Photonics  
solutions*

Mr. Kevin Tam  
*Usyd*

Dr. Chad Husko  
*Argonne Nat. Lab*

## Topological quantum work



Dr. Bryn Bell  
*Imperial College*

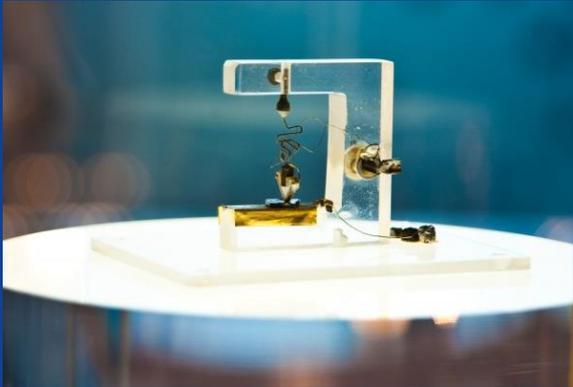
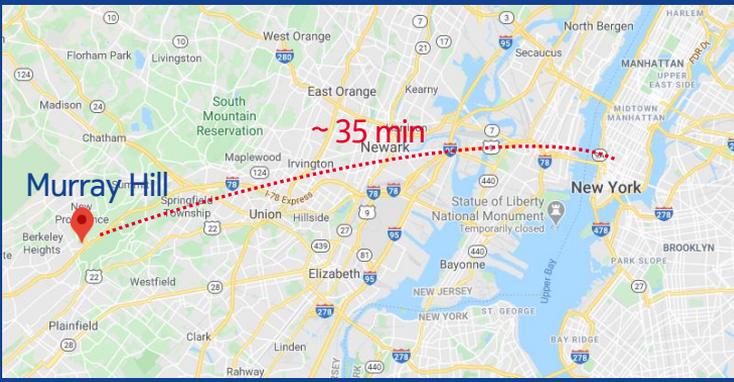
Ms. Michelle Wang  
*Usyd*

Mr. Cooper Doyle  
*Usyd*

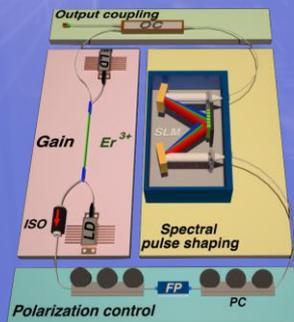
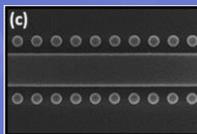
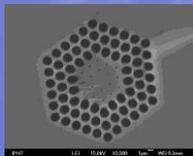
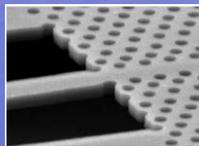
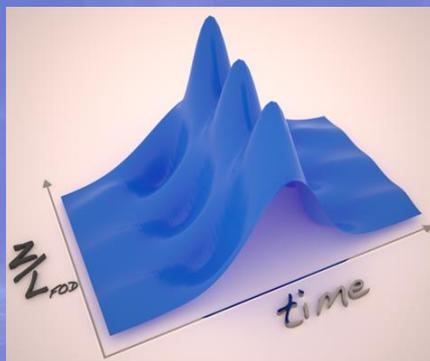
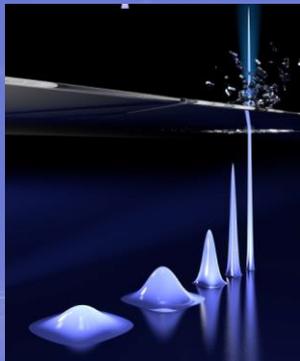
Prof. Ben Eggleton  
*Usyd*

Prof. Moti Segev  
*Technion*

# Murray Hill, New Providence, NJ – BELL LABS HEADQUARTERS

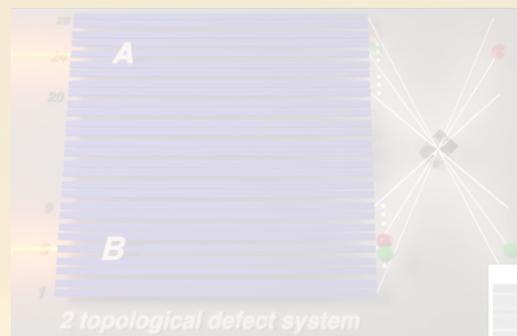


## PART 1: NOVEL SOLITON PHYSICS AND TECHNOLOGIES



Runge et al. *Phys. Rev. Res.* **3**, 013166 (2021)  
 Runge et al. *Nat. Photonics* **14**, 492-497 (2020).  
 Sahin, Blanco-Redondo, et al. *Laser & Photonics Reviews* **13**, (2019)  
 Blanco-Redondo et al. *Nature Comm.* **7**, 10427 (2016)  
 Blanco-Redondo et al. *Nature Comm.* **5**, 3160 (2014)  
 Blanco-Redondo et al. *Optica* **1**, 299-306 (2014)

## PART 2: TOPOLOGICAL QUANTUM PHOTONICS



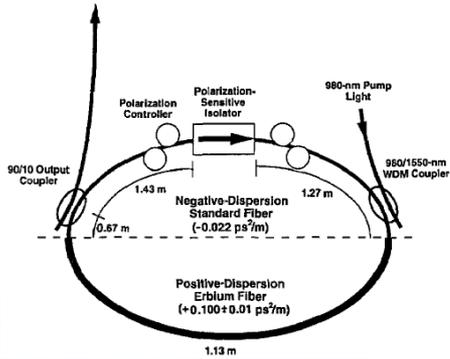
Doyle et al. (in preparation)  
 Khan et al. *Advanced Materials Technologies*, 2100252 (2021)  
 Blanco-Redondo, *Proceedings of the IEEE* **108**, 837-849 (2020)  
 Wang et al. *Nanophotonics* **8**, 1327-1335 (2019)  
 Blanco-Redondo et al. *Science* **362**, 568-571 (2018)  
 Blanco-Redondo et al. *Phys. Rev. Lett.* **116**, 163901 (2016)

# Solitons: the interplay of negative dispersion and positive nonlinearity

## Other cavity configurations :

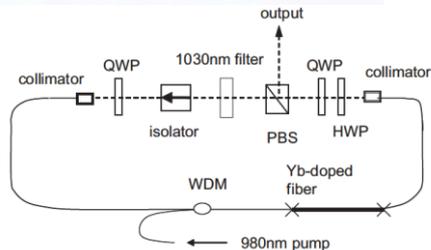
### Stretched-pulse laser invented @MIT

Tamura, Ippen, Haus & Nelson, Opt. Lett. (1993)



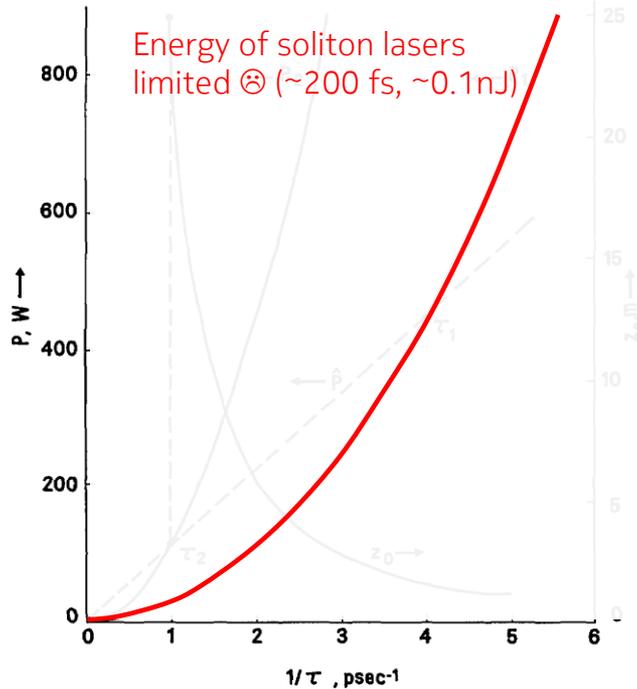
### ANDi laser invented @Cornell

Chong, Buckley, Rensinger & Wise, Opt. Exp. (2006)



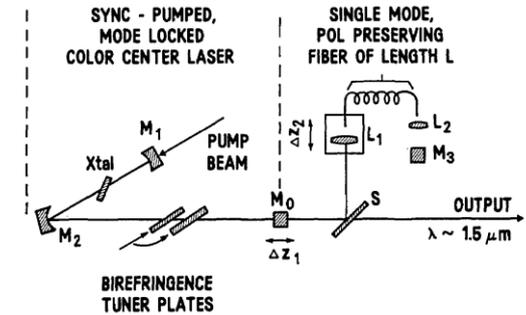
Solitons in optical fiber predicted & observed @Bell Labs  
 Mollenauer & Tappert, APL (1973)  
 Mollenauer, Stolen, & Gordon, Opt. Lett. (1976)

Energy of soliton lasers limited ☹ (~200 fs, ~0.1nJ)



### The soliton laser invented @Bell Labs

Mollenauer & Stolen Opt. Lett. (1984)



“... the soliton laser would most probably take the form of a single loop of fiber closed upon itself. The simplicity and low cost of such devices would make them most attractive.”

Fixed Energy-width relation:  
 a blessing and a curse

$$E_S \propto \frac{1}{\tau}$$

Well-defined shape (sech²)

Transformed limited pulses (unchirped)

Bell Labs

L. Mollenauer, R. Stolen, J. Gordon

# Soliton functionality in silicon

A complicated affair

Negative  $\beta_2$  required

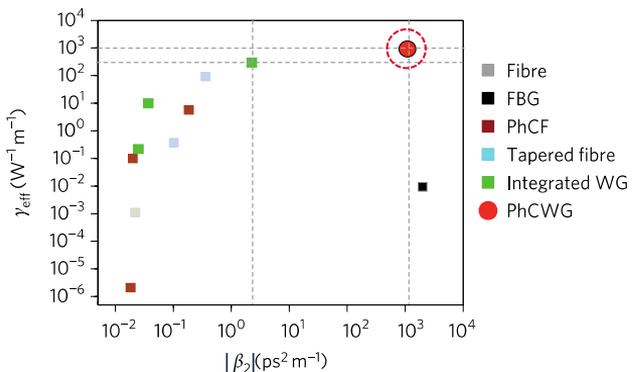
Nonlinear Schrödinger Eq.

$$\frac{\partial A}{\partial z} = -i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial T^2} + i \gamma |A|^2 A$$

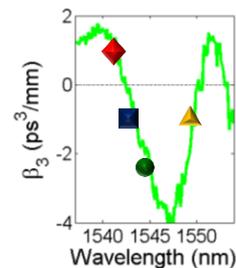
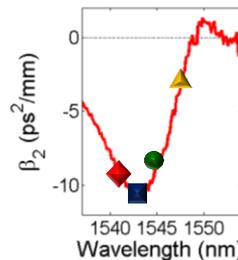
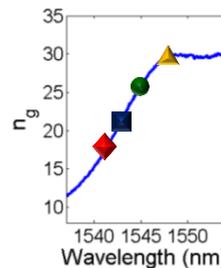
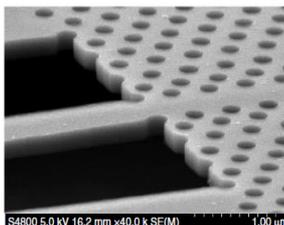
Negative  $\beta_2$  required

Generalized Nonlinear Schrödinger Eq. in Silicon

$$\frac{\partial A}{\partial z} = -\frac{\alpha_l}{2} A - i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial T^2} + \frac{\beta_3}{6} \frac{\partial^3 A}{\partial T^3} + \left( i \gamma - \frac{\gamma_{TPA}}{2} \right) |A|^2 A + \left( i k_0 n_{FC} - \frac{\sigma}{2} \right) N_c A$$



Colman\*, Husko\*, Combrie\* et al, *Nature Photonics* 4, 862 (2010)



▲ **Solitons:** moderate anomalous dispersion, high  $n_g$  (nonlinearity), low third order dispersion

Blanco-Redondo et al, *Nature Commun.* 5, 3160 (2014)

◆ ■ ● **Free-carriers dominate nonlinear pulse propagation:**

- ❖ Anomalous dispersion + Free-carrier dispersion → Nonlinear temporal broadening (x2)
- ❖ Third order dispersion can counteract this and other nonlinear free-carrier effects

Blanco-Redondo et al, *Optica* 1(5), 299 – 306 (2014)

# Pure-quartic solitons in silicon

A novel kind of pulse

Negative  $\beta_2$  required

Biharmonic Nonlinear Schrödinger Eq.

$$\frac{\partial A}{\partial z} = i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial T^2} + i g \|A\|^2 A$$

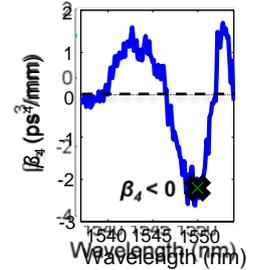
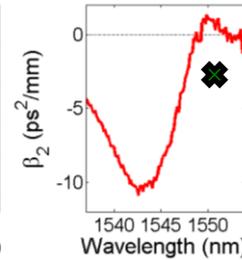
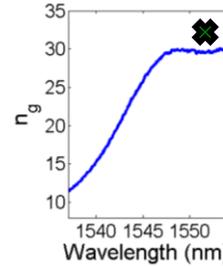
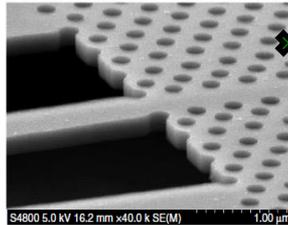
Negative  $\beta_4$  required

Negative  $\beta_2$  required

Generalized nonlinear Schrödinger Eq. in Silicon

$$\text{vs. } \frac{\partial A}{\partial z} = \frac{\partial A}{\partial z} \frac{\alpha_1}{2} A + \frac{\alpha_1}{2} A - i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial T^2} + i \frac{\beta_3}{2} \frac{\partial^3 A}{\partial T^3} + i \frac{\beta_4}{24} \frac{\partial^4 A}{\partial T^4} + \left( \frac{\gamma_{TPA}}{2} \|A\|^2 + \frac{\gamma_{FPA}}{2} \|A\|^4 \right) A + \left( A k_0^2 A_{FC} + \left( i \frac{\sigma}{2} \right) N_{FC} A - \frac{\sigma}{2} N_c A \right)$$

Negative  $\beta_4$  required



✱ Dominant (negative) quartic dispersion balanced by nonlinearity: **Pure-quartic solitons**

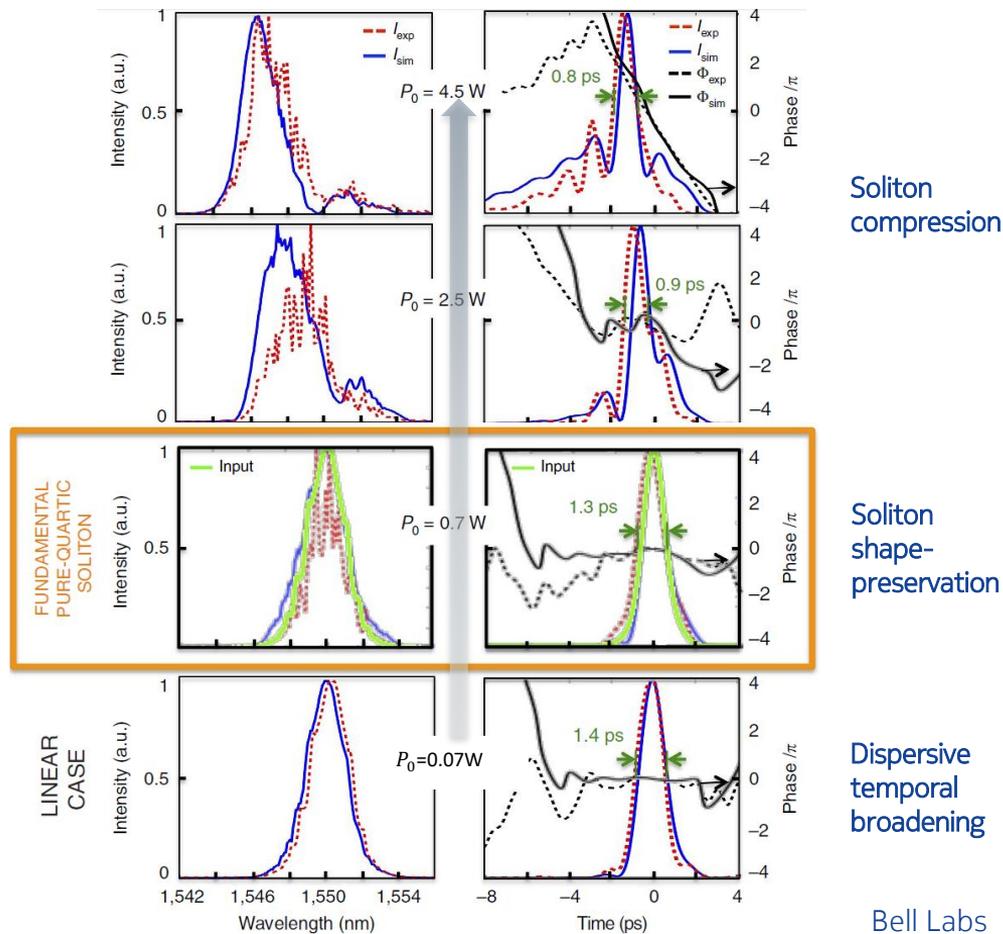
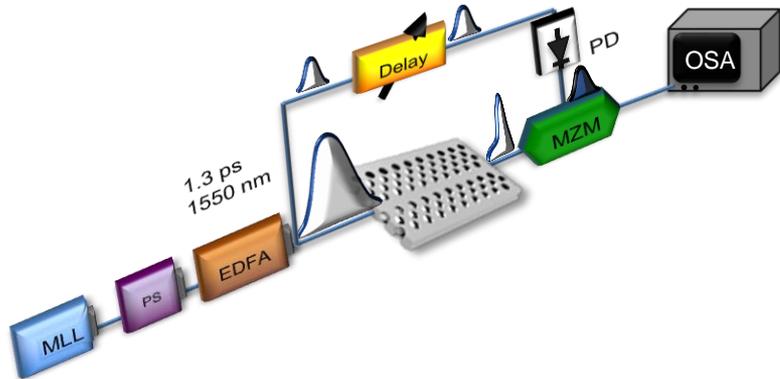
Blanco-Redondo et al, *Nature Commun.* 7, 10427 (2016)

# Pure-quartic solitons in silicon

A novel kind of pulse

Frequency resolved electrical gating

Dorrer and Kang, Opt. Lett 27(15), 1315-1317 (2002)



Soliton compression

Soliton shape-preservation

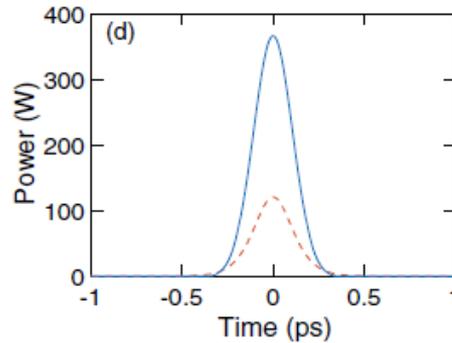
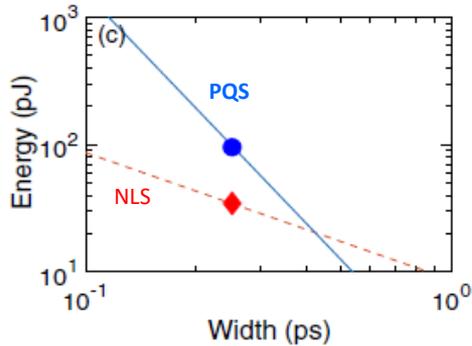
Dispersive temporal broadening

Bell Labs

Blanco-Redondo et al, *Nature Commun.* 7, 10427 (2016)

# Pure-quartic solitons

What is useful about them?



- Shape differs from *conventional* NLS soliton.
- Time-bandwidth product:

$$TBP_{PQS} = 0.53 \text{ vs } TBP_{NLS} = 0.315$$

- Different energy-width scaling relations!

$$E_{PQS} = \frac{2.87|\beta_4|}{\gamma\tau^3}$$

$$E_{PQS} \propto \frac{1}{\tau^3} \text{ vs } E_{NLS} \propto \frac{1}{\tau}$$

K.K.K. Tam, T. Alexander, A. Blanco-Redondo and C. Martijn de Sterke *Opt. Lett.* **44**, 3306 (2019).

A PQS laser could emit transform-limited pulses with higher energies than conventional NLS soliton lasers at short durations

# The pure-quartic soliton laser

## Overcoming energy limitations in soliton lasers

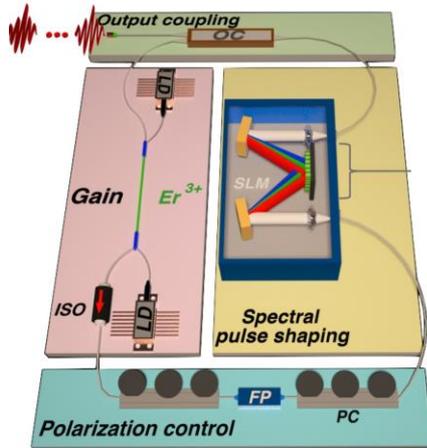
Need a laser cavity with **dominant negative quartic dispersion**

- How about photonic crystal fibers? Yes, but challenging

Lo, Stefani, Martijn de Sterke, and Blanco-Redondo *Opt. Exp.* **26**(6), 7786-7796 (2018)



- How about a dispersion-managed approach using a pulse-shaper?



- Soliton laser using NPE for mode-locking.
- **Intracavity pulse-shaper** allows for the tuning of the **net cavity dispersion**.

J. Schröder, et al. *Opt. Express.* **18**, 22715 (2010).

J. Peng and S. Boscolo *Sci. Rep.* **6**, 25995 (2016).

A. J. Runge, D. D. Hudson, K. K. K. Tam, C. Martijn de Sterke, and A. Blanco-Redondo, "The pure-quartic soliton laser," *Nature Photonics* **14**, 492-497 (2020)

# The pure-quartic soliton laser

## Overcoming energy limitations in soliton lasers

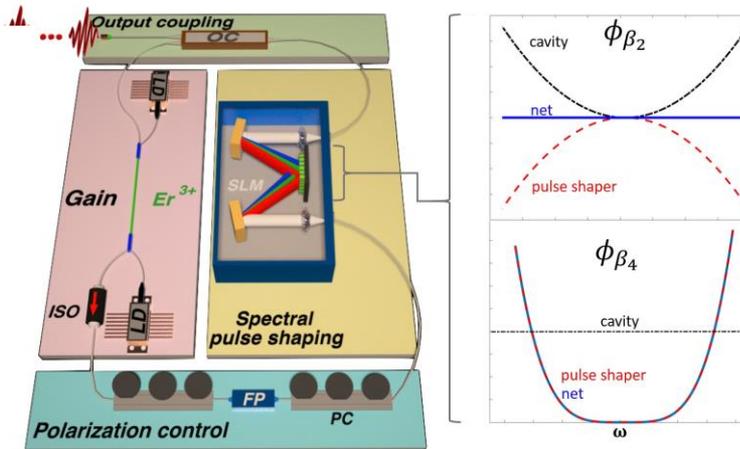
Need a laser cavity with **dominant negative quartic dispersion**

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Lo, Stefani, Martijn de Sterke, and Blanco-Redondo *Opt. Exp.* 26(6), 7786-7796 (2018)

- How about a dispersion-managed approach using a pulse-shaper?



- Pulse-shaper applies a **phase mask** to compensate for  $\beta_2$  and  $\beta_3$  of the fibres, and induces a **large negative quartic** ( $\beta_4$ ) dispersion.
- Cavity length  $L = 21.4$  m, opposite 2<sup>nd</sup> and 3<sup>rd</sup> order dispersion of SMF  $\beta_2 = +21.4$  ps<sup>2</sup>/km,  $\beta_3 = -0.2$  ps<sup>3</sup>/km.
- $\beta_4 = -80$  ps<sup>4</sup>/km is chosen to be much larger than intrinsic 4<sup>th</sup> order dispersion of SMF

A. J. Runge, D. D. Hudson, K. K. K. Tam, C. Martijn de Sterke, and A. Blanco-Redondo, "The pure-quartic soliton laser," *Nature Photonics* 14, 492-497 (2020)

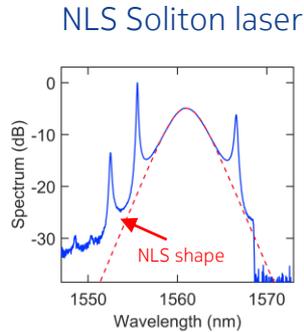
# The pure-quartic soliton laser

## Overcoming energy limitations in soliton lasers

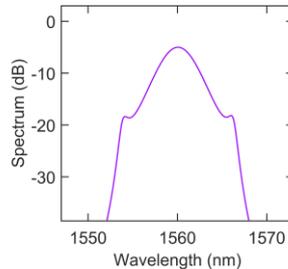


Measurements by  
Dr. Antoine Runge

MEASUREMENTS



LASER  
SIMULATIONS

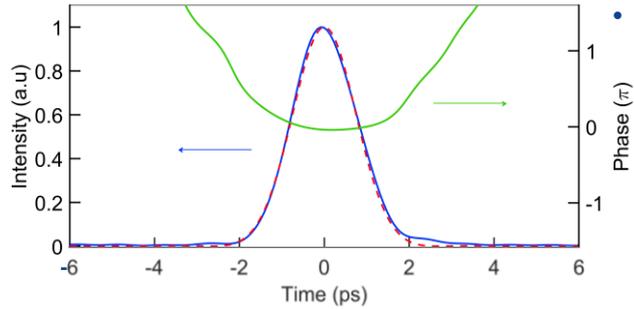


- Phase mask **OFF**, laser operates in conventional soliton regime.
- Phase mask **ON**, different output pulse spectrum (blue). Good fit with calculated PQS shape (dashed red).

A. J. Runge, D. D. Hudson, K. K. K. Tam, C. Martijn de Sterke, and A. Blanco-Redondo, "The pure-quartic soliton laser," *Nature Photonics* 14, 492-497 (2020)

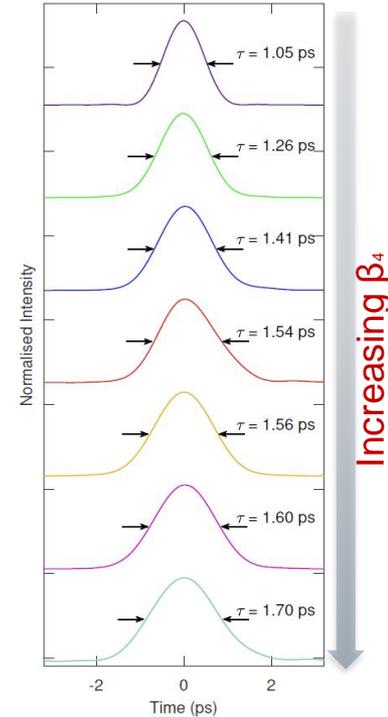
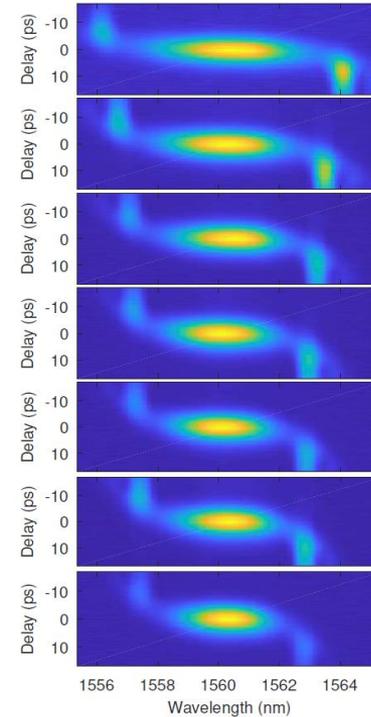
# The pure-quartic soliton laser

## Overcoming energy limitations in soliton lasers



- Nearly-flat phase  $\rightarrow$  Nearly transform limited

- Adjustable spectrum and temporal profile by changing the dispersion provided to the pulse shaper

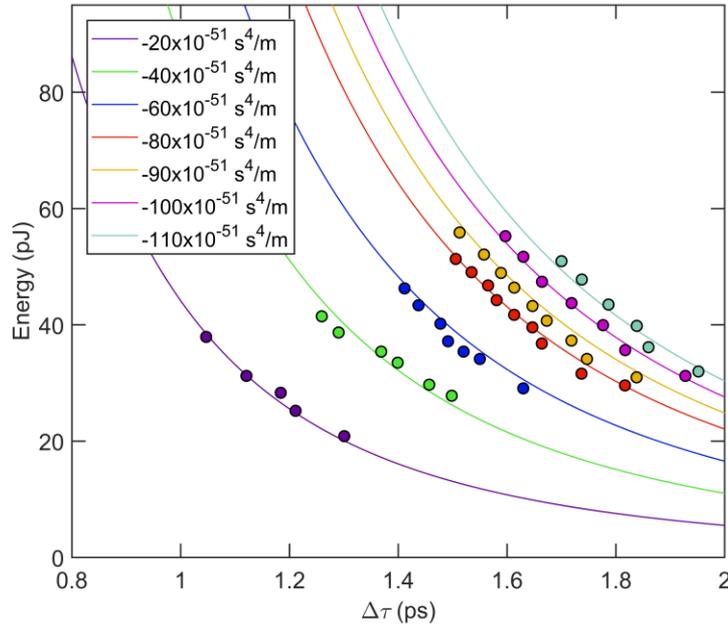


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# The pure-quartic soliton laser

## Overcoming energy limitations in soliton lasers

### Energy-width scaling relation

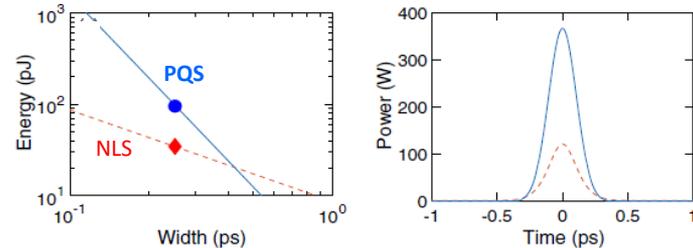


- We have demonstrated that the energy scales as predicted:

$$E_{PQS} = \frac{2.87|\beta_4|}{\gamma\tau^3}$$

- First step towards solving soliton laser limitations:

$$E_{PQS} \propto \frac{1}{\tau^3} \text{ vs } E_{NLS} \propto \frac{1}{\tau}$$

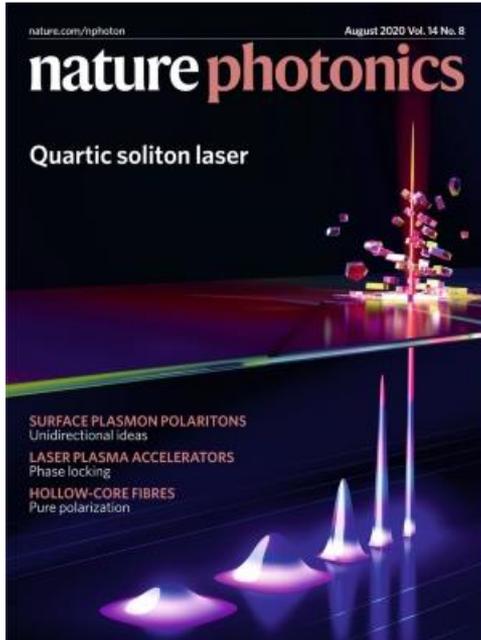


*“The simplicity and low cost of such devices would make them most attractive.”*  
Mollenauer and Stolen, Opt. Lett. (1984)

A. J. Runge, D. D. Hudson, K. K. K. Tam, C. Martijn de Sterke, and A. Blanco-Redondo, “The pure-quartic soliton laser,” *Nature Photonics* 14, 492-497 (2020)

# The pure-quartic soliton laser

## Overcoming energy limitations in soliton lasers



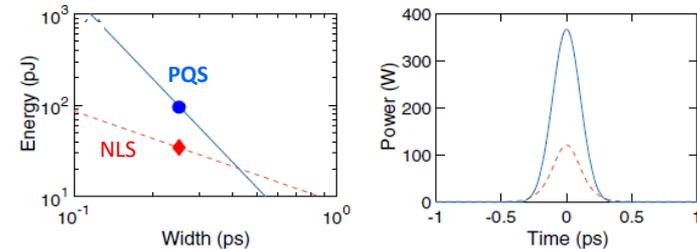
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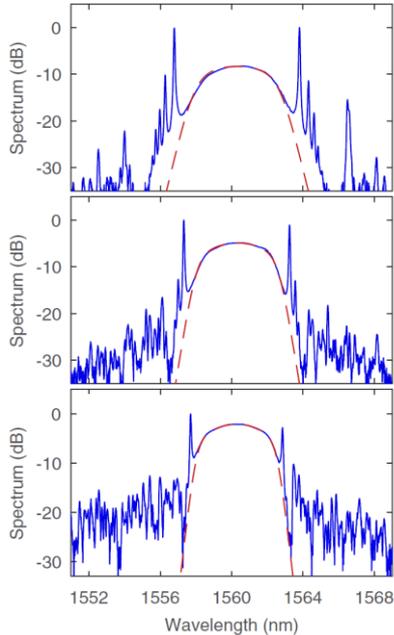
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A. J. Runge, D. D. Hudson, K. K. K. Tam, C. Martijn de Sterke, and A. Blanco-Redondo, “The pure-quartic soliton laser,” *Nature Photonics* (2020) <https://doi.org/10.1038/s41566-020-0629-6>

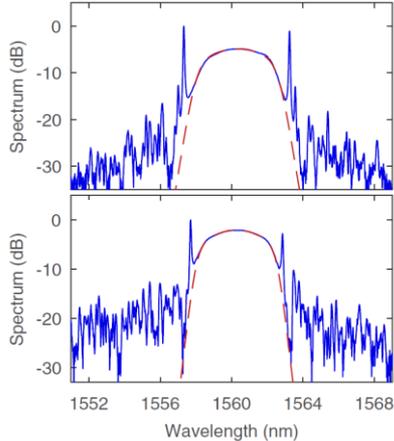
# How about other even orders of dispersion?

Pure-sextic, pure-octic, pure-decic soliton lasers?

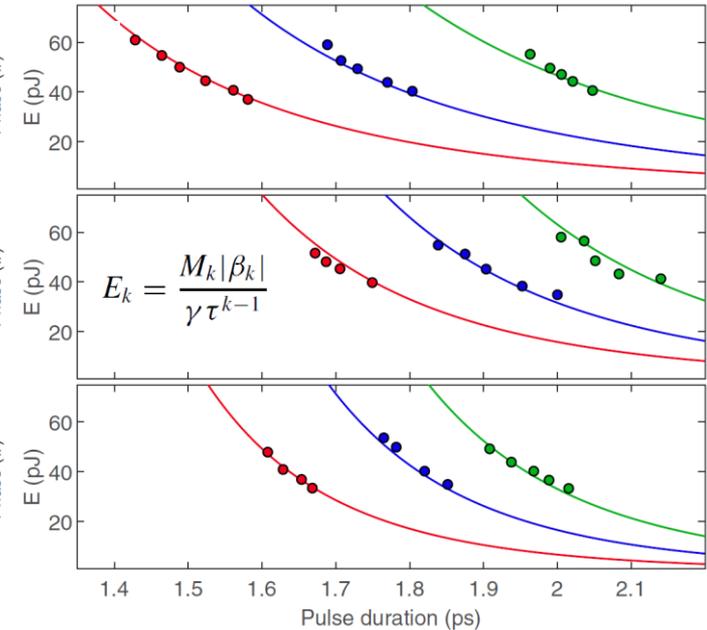
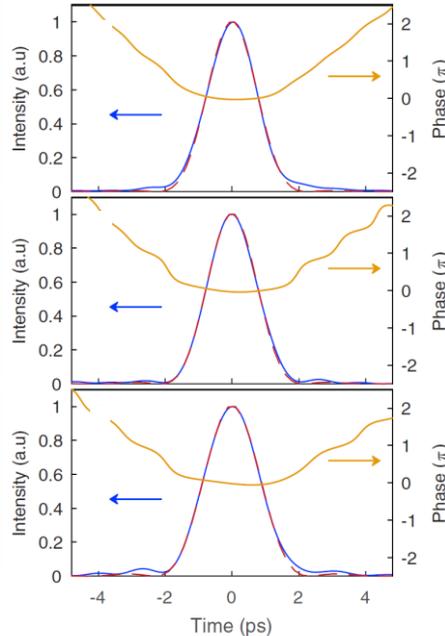
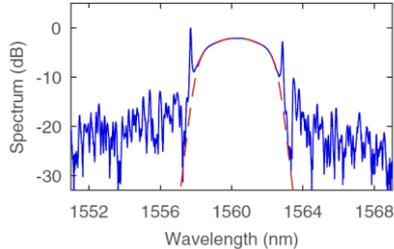
Pure-sextic soliton  
 $\beta_6$



Pure-octic soliton  
 $\beta_8$



Pure-decic soliton  
 $\beta_{10}$



Runge, Hudson, Martijn de Sterke and Blanco-Redondo CLEO 2020 (paper JTh4B.1) - Postdeadline Paper  
Runge, Qiang, Alexander, Rafat, Hudson, Blanco-Redondo and Martijn-de Sterke, Phys. Rev. Res. 3, 013166 (2021)

# End of Part 1

## Conclusions

Solitons and other nonlinear effects in silicon are challenging but possible

Demonstrated new type of mode-locked laser emitting pure-quartic, pure-sixtic, pure-octic... solitons

- Energy-width scaling

The desired higher-order dispersion profile can be achieved via:

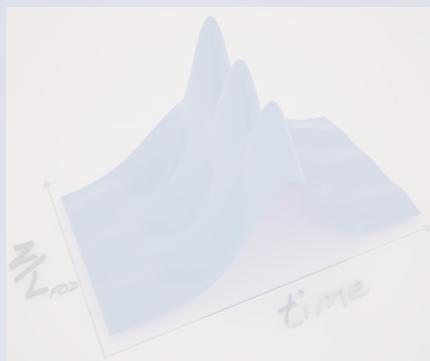
- Dispersion engineering
- Pulse shaping element

Higher-order dispersion can be more  
than the ugly duckling



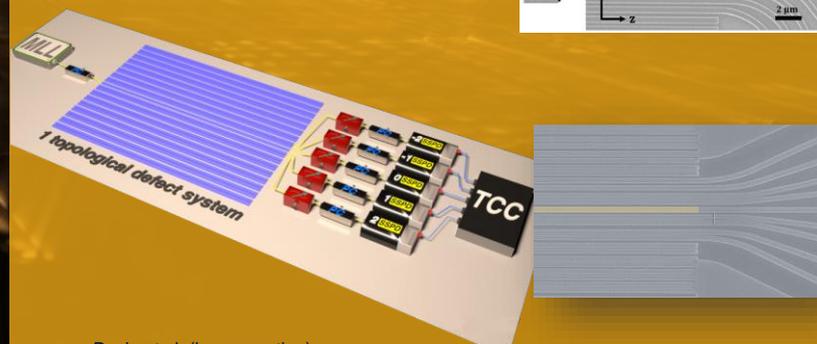
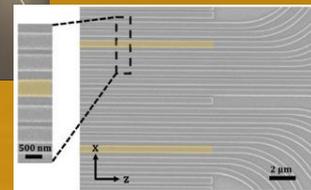
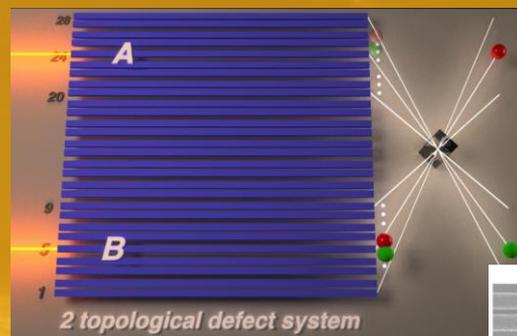
Bell Labs

## PART 1: NOVEL SOLITON PHYSICS AND TECHNOLOGIES



Runge et al. *Nature Photonics* *Nat. Photonics* (2020). <https://doi.org/10.1038/s41566-020-0629-6>  
 Tam et al. *Physical Review A* **101**, 043822 (2020)  
 Tam et al. *Opt. Lett.* **44**, 3306 (2019)  
 Sahin et al. *Laser & Photonics Reviews* **13**(8) (2019)  
 Blanco-Redondo et al. *Nature Comm.* **7**, 10427 (2016)  
 Blanco-Redondo et al. *Nature Comm.* **5**, 3160 (2014)  
 Blanco-Redondo et al. *Optica* **1**, 299-306 (2014)

## PART 2: TOPOLOGICAL QUANTUM PHOTONICS



Doyle et al. (in preparation)  
 Khan et al. *Advanced Materials Technologies*, 2100252 (2021)  
 Blanco-Redondo, *Proceedings of the IEEE* **108**, 837-849 (2020)  
 Wang et al. *Nanophotonics* **8**, 1327-1335 (2019)  
 Blanco-Redondo et al. *Science* **362**, 568-571 (2018)  
 Blanco-Redondo et al. *Phys. Rev. Lett.* **116**, 163901 (2016)

# What is topology?

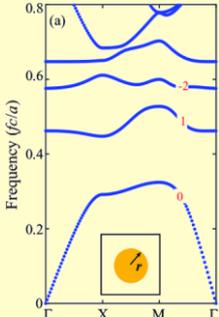
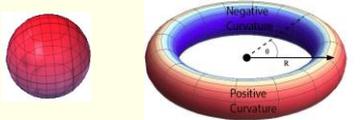
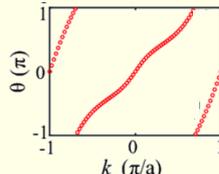
Topologies – quantities that DO NOT CHANGE when we CONTINUOUSLY deform a geometric object



Number of holes preserved under continuous deformation

**# holes = genus = topology**

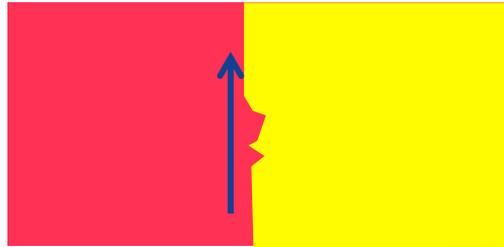
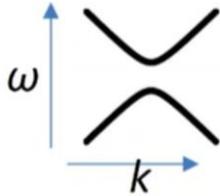
# Topology in geometric objects (real space) vs in photonic systems (momentum space)

Real Space	Momentum space
<p>Geometric objects</p> 	<p>Photonic bands</p> 
<p>Gaussian curvature</p> $\mathcal{K}(r) = 1/R^2$ 	<p>Berry curvature</p> $\mathcal{F}(\mathbf{k}) = \nabla_{\mathbf{k}} \times \mathcal{A}(\mathbf{k})$ 
<p>Genus (Number of holes)</p> $\frac{1}{2\pi} \int_{\text{surface}} \mathcal{K} dA = 2(1 - g)$ 	<p>Chern number (Number of twists/untwists)</p> $C = \frac{1}{2\pi} \oint \mathcal{F}(\mathbf{k}) \cdot d\mathbf{s}$ 

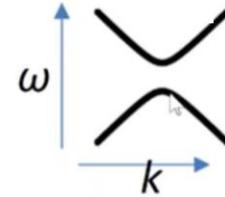
## Topological edge states in photonic systems

- ✓ Topological invariant (Chern number or Zak Phase) can only take integer values
- ✓ Topological invariant of gapped systems (insulators) cannot change without closing the gap  
→ Topological phase transition

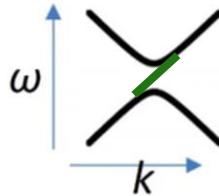
$$C = \frac{1}{2\pi} \oint \mathcal{F}(\mathbf{k}) \cdot d\mathbf{s} = 0$$



$$C = \frac{1}{2\pi} \oint \mathcal{F}(\mathbf{k}) \cdot d\mathbf{s} = 1$$



**Unidirectional, gapless states at the interface between two materials with different topological invariants**



# Our topological platform:

## A silicon photonics implementation of the SSH model

Blanco-Redondo, et al., *Phys. Rev. Lett.* **116**, 163901 (2016)

### Other SSH implementations:

Su, Schrieffer, and Heeger, PRL (1979)

Malkova et al. Opt. Lett. (2009)

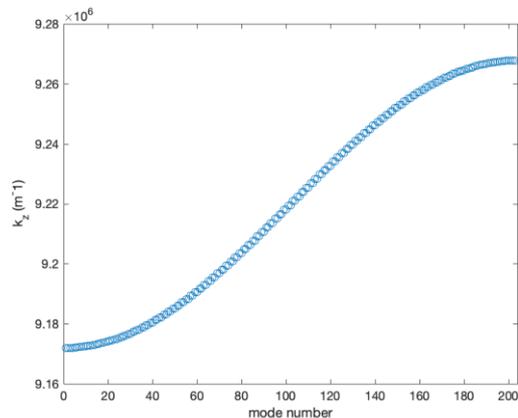
Zeuner et al. PRL (2015)

Poli et al. Nat. Comm. (2015)

### Relevant topological invariant – Zak phase:

$$\mathcal{Z} = i \oint dq \langle u_q | \partial_q u_q \rangle$$

Zak, PRL 62, 2747 (1989)

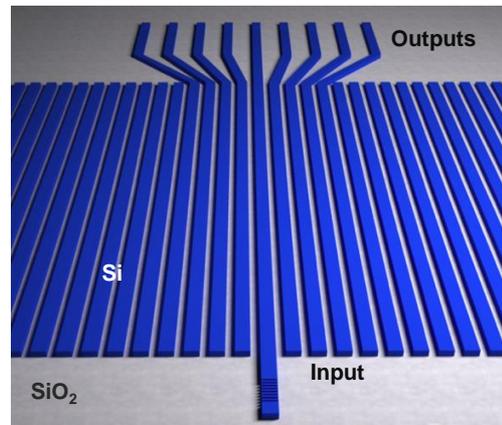
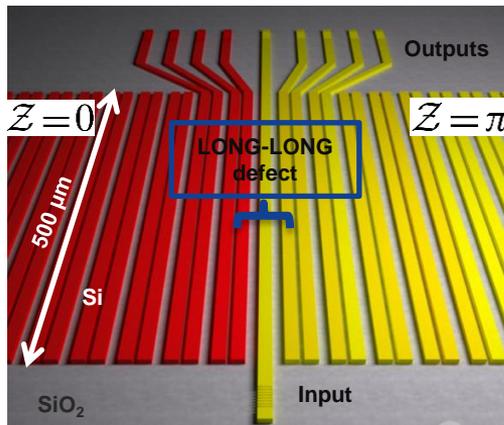


### Si waveguides:

- width:  $w = 450$  nm
- height:  $h = 220$  nm
- length:  $L = 500$   $\mu\text{m}$

### Gaps:

- 166 nm
- 294 nm

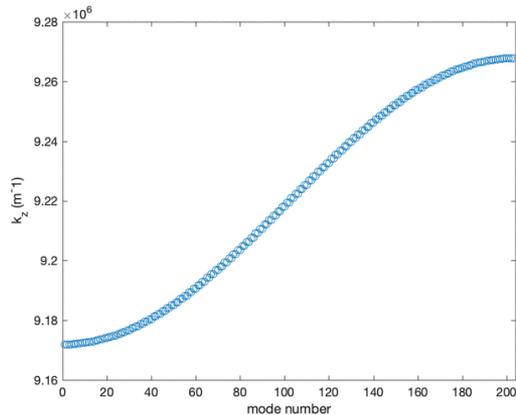


# Our topological platform:

## A silicon photonics implementation of the SSH model

Blanco-Redondo, et al., *Phys. Rev. Lett.* **116**, 163901 (2016)

From equidistant to SSH: ■ ■ ■ ■ ■ ■ ■ ■ ■ ■



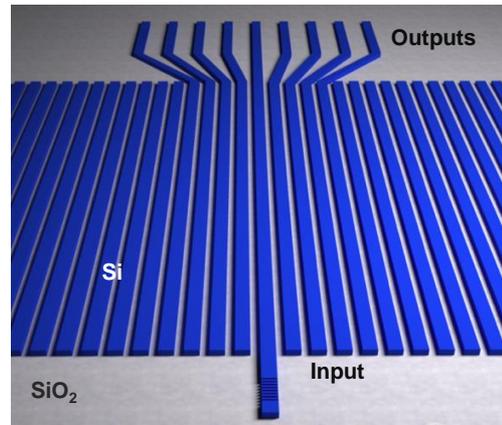
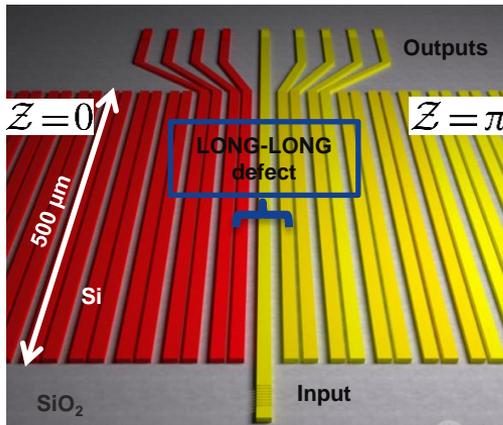
Blanco-Redondo, et al., *Phys. Rev. Lett.* **116**, 163901 (2016)

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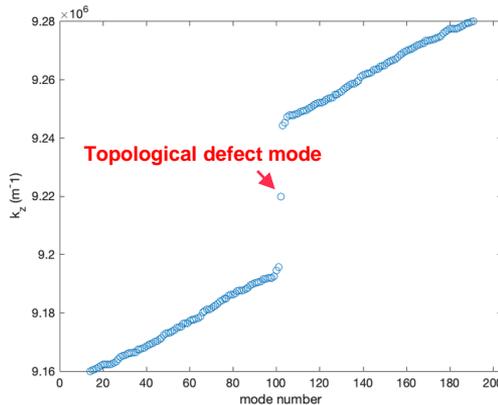


# Our topological platform:

## A silicon photonics implementation of the SSH model

Blanco-Redondo, et al., *Phys. Rev. Lett.* **116**, 163901 (2016)

Random disorder in waveguides positions:

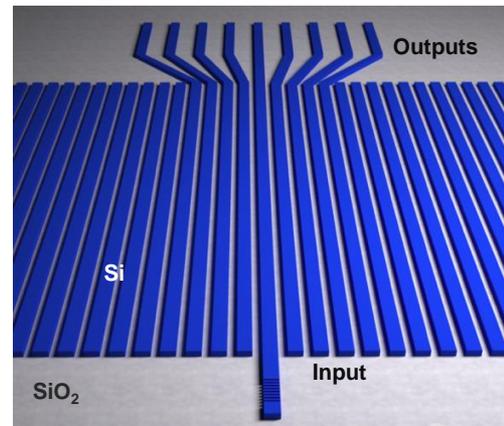
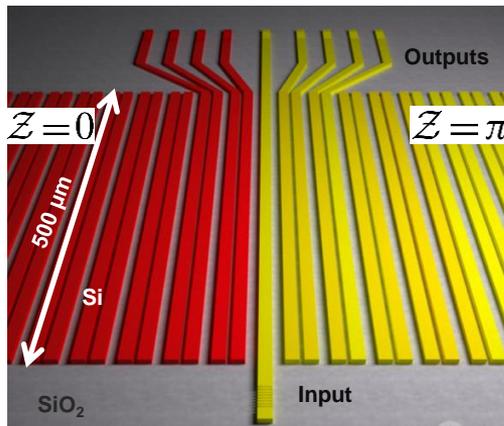


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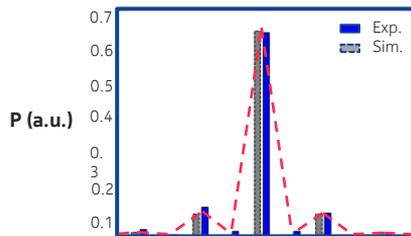


# Our topological platform:

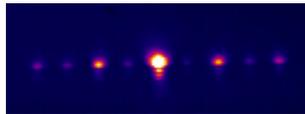
## A silicon photonics implementation of the SSH model

Blanco-Redondo, et al., *Phys. Rev. Lett.* **116**, 163901 (2016)

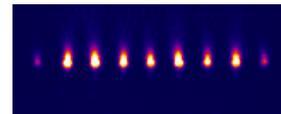
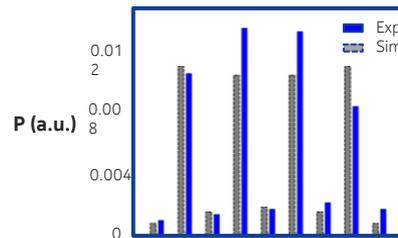
Topological localization



NIR images



No localization - Discrete diffraction

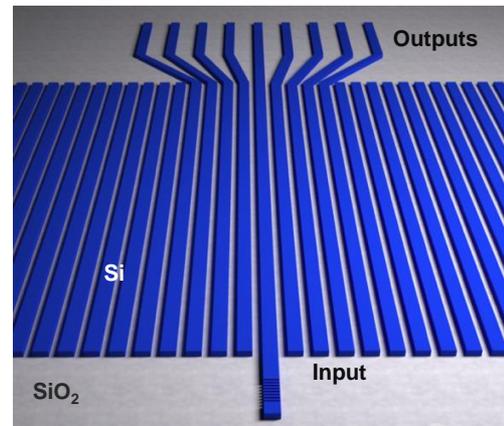
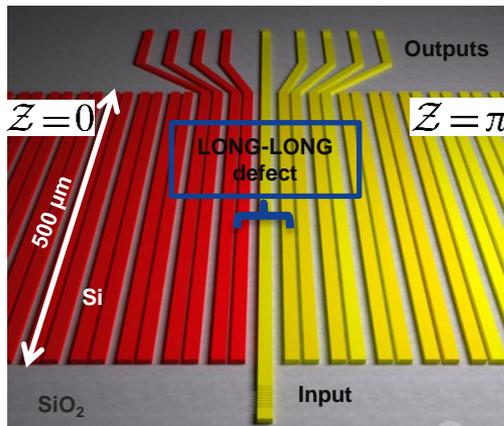


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**Gaps:**

- 166 nm
- 294 nm



# Topological quantum experiments on-chip

Can topology protect photonic quantum states?

Ref.	Geometry	Material	Model	Generation of quantum state	Quantum state	Topologically protected features	
Barik et al. <i>Science</i> 359,666-668, 2018	2D photonic crystals	GaAs	QSH	InGaAs quantum dots	Single photon	Unidirectional single-photon transport	M. Hafezi's group (JQI, Maryland)
Mittal et al. <i>Nature</i> 561, 502-506, 2018	2D coupled resonator array	Si	QSH	SFWM on-chip	Correlated photon pairs	Spectral correlations	
Blanco-Redondo et al. <i>Science</i> 372, 568-571, 2018	1D waveguide (wg) array	Si	SSH	SFWM on-chip	Correlated photon pairs	Spatial correlation biphoton	<b>Our work</b>
M. Wang et al. <i>Nanophotonics</i> 8, 1327-1335, 2019	1D wg. array	Si	SSH	SFWM on-chip	Path-entangled states	Path-entangled biphoton correlation	
Tambasco et al. <i>Science Adv.</i> 4, eaat3187, 2018	1D wg. array	Borosilicate	AAH	SPDC off-chip	Heralded single photon	High-visibility quantum interference	A. Peruzzo's group (MIT)
Y. Wang et al. <i>PRL</i> 122, 193903, 2019	1D wg. array	Borosilicate	SSH	SPDC off-chip	Heralded single photon	Second-order anti-correlation	Xianmin Jin's group (SJTU, Shanghai)
Y. Wang et al. <i>Optica</i> 6, 955-960, 2019	1D wg. array	Borosilicate	AAH	SPDC off-chip	Correlated photon pairs	Cross-correlation function	
Y. Wang et al. arXiv:1903.03015	1D wg. array	Borosilicate	SSH	SPDC off-chip	Polarization entangled states	High-concurrence and high-purity	

**Review paper:** A. Blanco-Redondo, "Topological Nanophotonics: Toward Robust Quantum Circuits," in *Proceedings of the IEEE* 108 (5), pp. 837-849, 2020

# Topological quantum experiments on-chip

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Mittal et al. <i>Nature</i> 561, 502-506, 2018	2D coupled resonator array	Si	QSH	SFWM on-chip	Single photon	Spectral correlations
Blanco-Redondo et al. <i>Science</i> 372, 568-571, 2018	1D waveguide (wg) array	Si	SSH	SFWM on-chip	Photon pairs	Spatial correlation
M. Wang et al. <i>Nanophotonics</i> 8, 1327-1335, 2019	1D wg. array	Si	SSH	SFWM on-chip	Photon pairs	Path-entangled biphoton correlation
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M. Hafezi's group (JQI, Maryland)

**Our work**

A. Peruzzo's group (RMIT)

Xianmin Jin's group (SJTU, Shanghai)

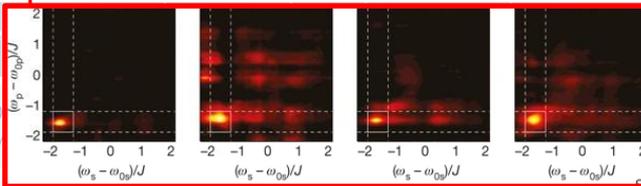
**Review paper:** A. Blanco-Redondo, "Topological Nanophotonics: Toward Robust Quantum Circuits," in *Proceedings of the IEEE* 108 (5), pp. 837-849, 2020

# Topological quantum experiments on-chip

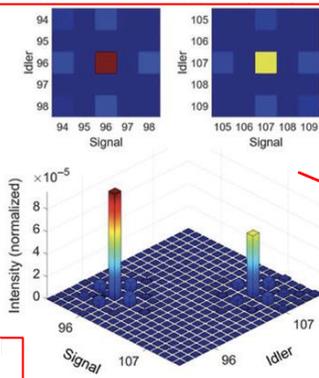
Can topology protect photonic quantum states?

Ref.	Ge	Si	SS	AA	Other	Features	Group
Barik et al. <i>Science</i> 359,666-668, 2018	2D cry					Unidirectional single-photon transport	M. Hafezi's group (JQI, Maryland)
Mittal et al. <i>Nature</i> 561, 502-506, 2018	2D coupled resonator array					Spectral correlations	M. Hafezi's group (JQI, Maryland)
Blanco-Redondo et al. <i>Science</i> 372, 568-571, 2018	1D waveguide (wg) array	Si	SS			Spatial correlation	<b>Our work</b>
M. Wang et al. <i>Nanophotonics</i> 8, 1327-1335, 2019	1D wg. array	Si	SS			Path-entangled biphoton correlation	
Tambasco et al. <i>Science Adv.</i> 4, eaat3187, 2018	1D wg. array	Borosilicate	AA			High-visibility quantum interference	A. Peruzzo's group (RMIT)
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Y. Wang et al. arXiv:1903.03015	1D					High-concurrence and high-purity	

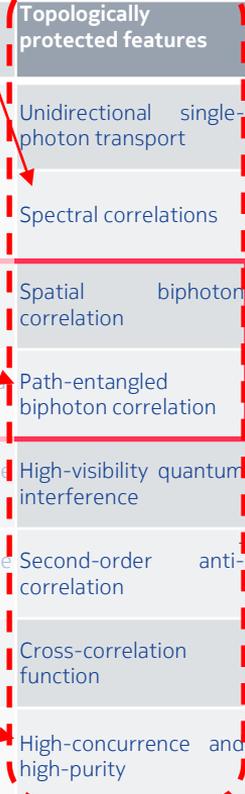
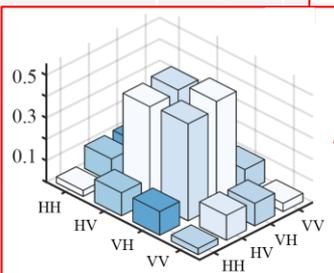
**Spectral correlations**



**Path entanglement**

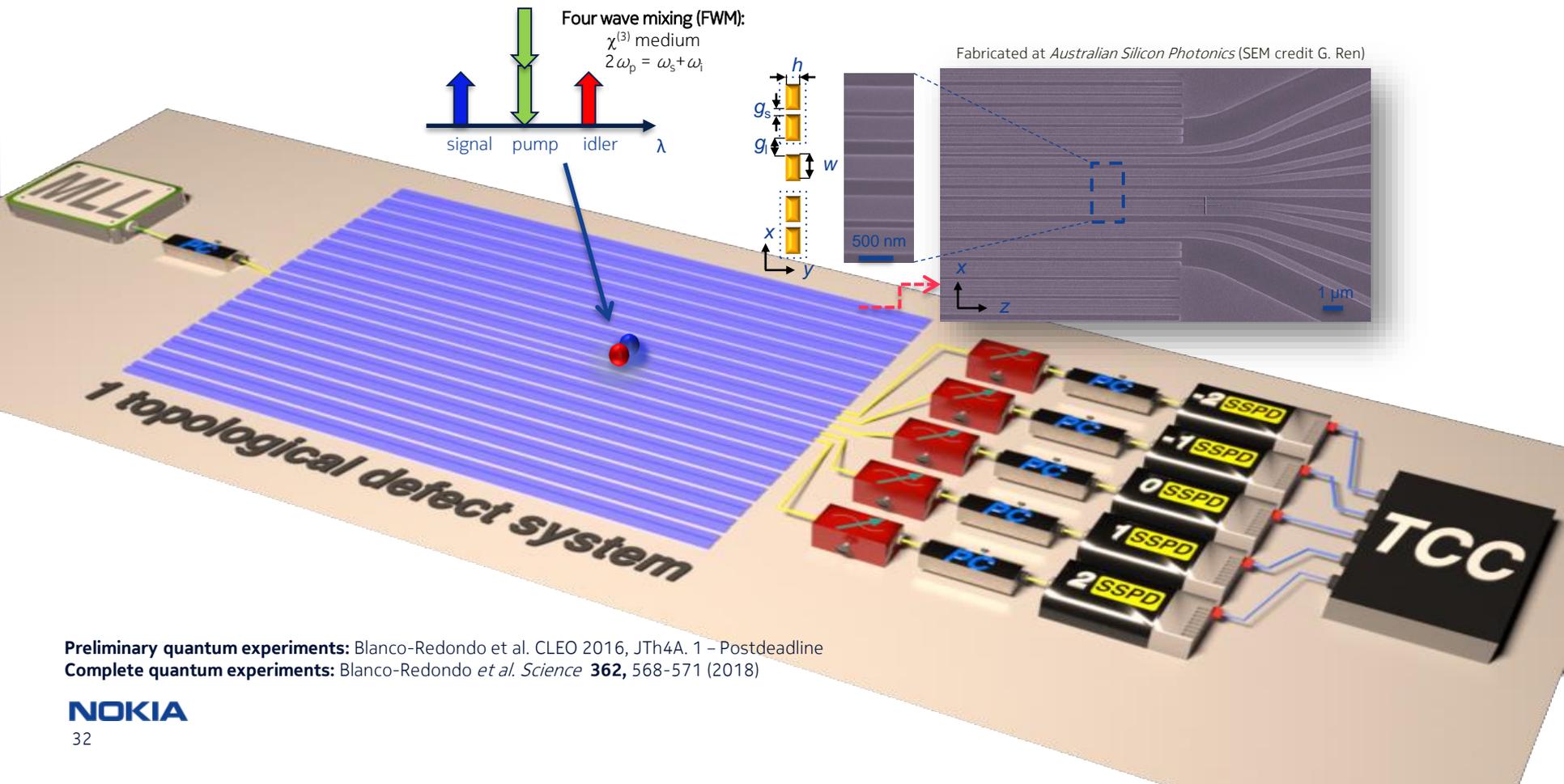


**Polarization entanglement**



**Review paper:** A. Blanco-Redondo, "Topological Nanophotonics: Toward Robust Quantum Circuits," in *Proceedings of the IEEE* 108 (5), pp. 837-849, 2020

# Topological protection of biphoton states

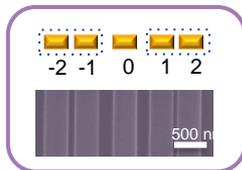


**Preliminary quantum experiments:** Blanco-Redondo et al. CLEO 2016, JTh4A. 1 – Postdeadline

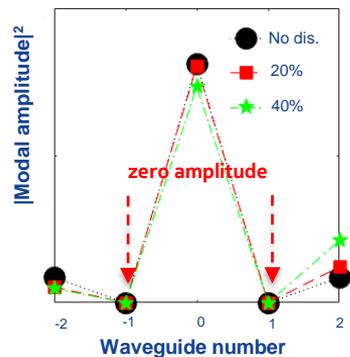
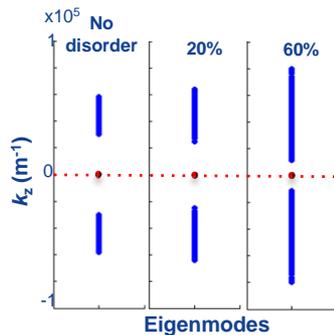
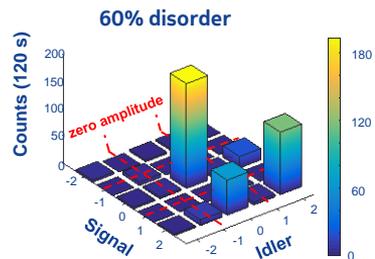
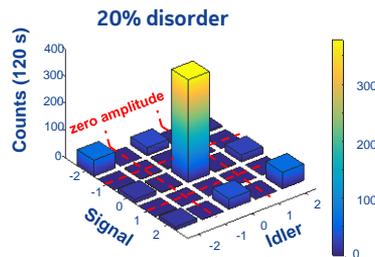
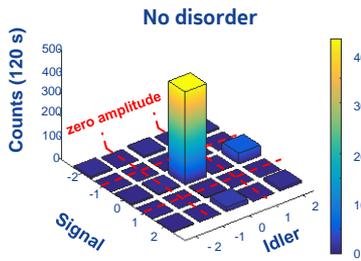
**Complete quantum experiments:** Blanco-Redondo *et al. Science* **362**, 568-571 (2018)

# Measuring the biphoton correlation at the output of our topological lattice

TOPOLOGICAL  
lattice

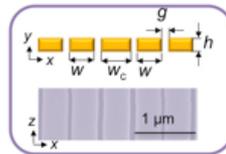


Blanco-Redondo, Bell, Oren, Eggleton & Segev, *Science* 362, 568-571 (2018)

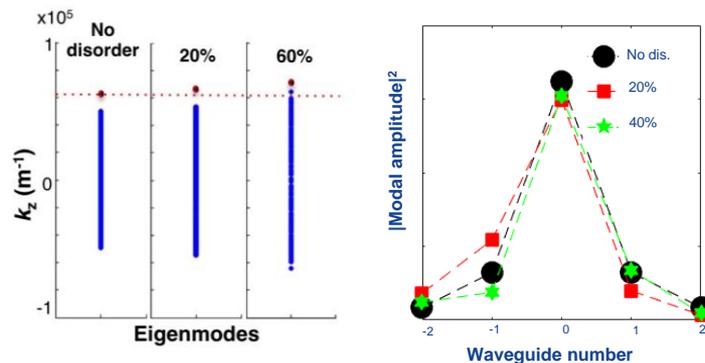
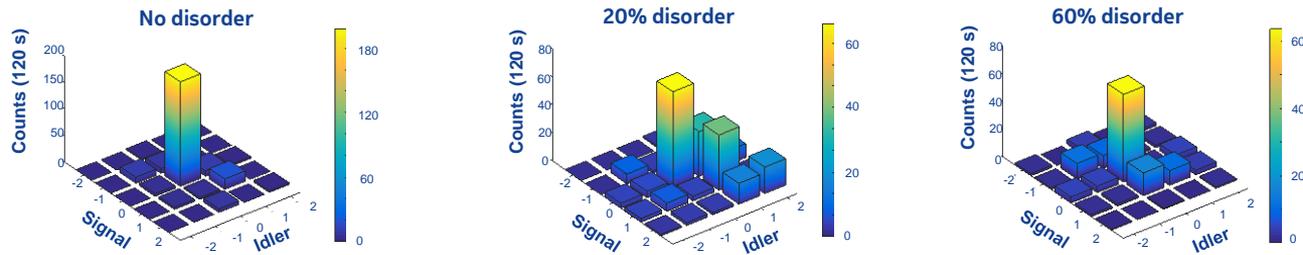


# Measuring the biphoton correlation at the output of our trivial lattice

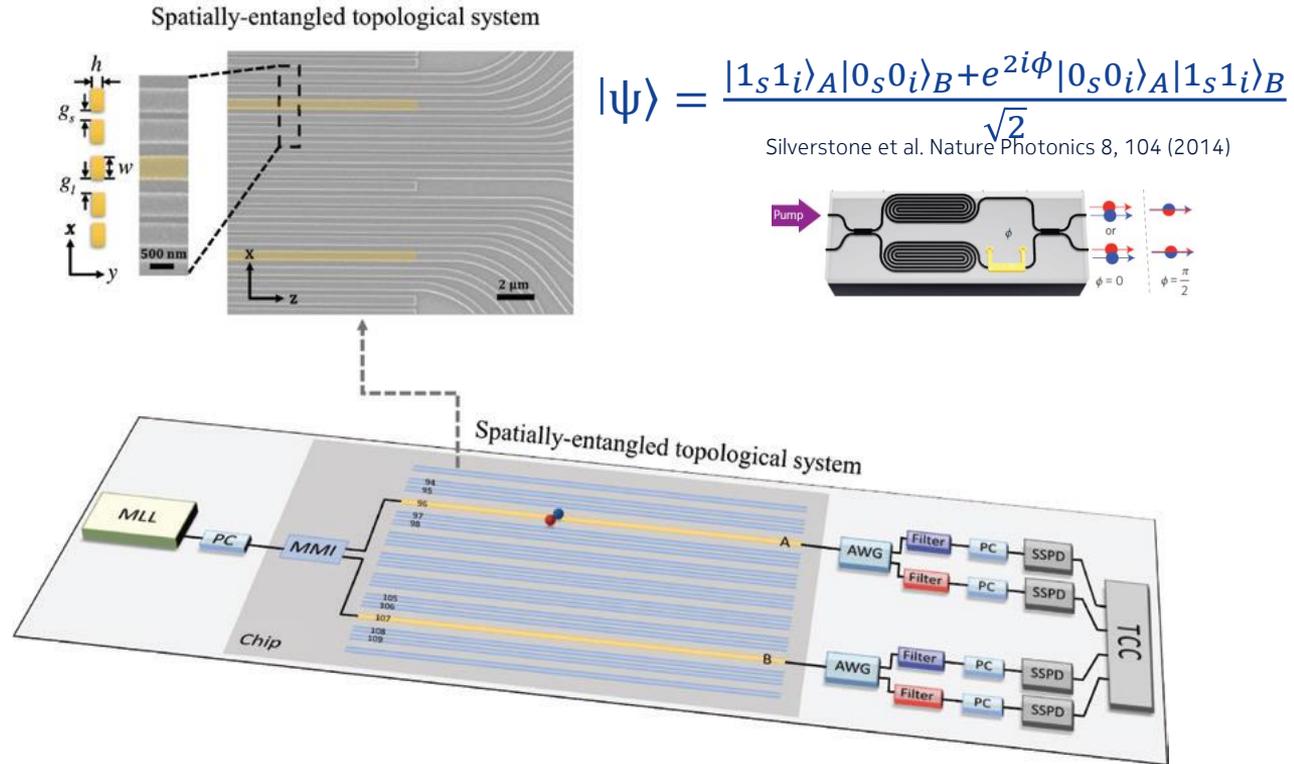
TRIVIAL  
lattice



Blanco-Redondo, Bell, Oren, Eggleton & Segev, *Science* 362, 568-571 (2018)

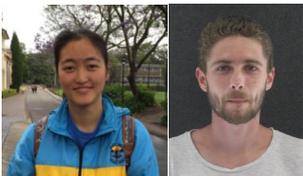


# Path-entanglement between topological modes



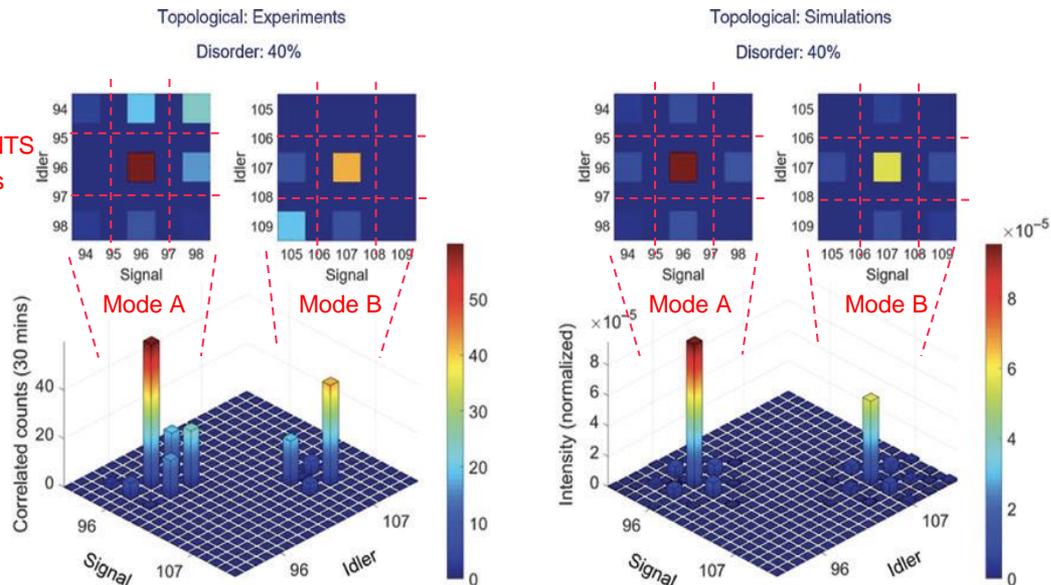
M. Wang, C. Doyle, B. Bell, M. J. Collins, E. Magi, B. J. Eggleton, M. Segev, and A. Blanco-Redondo, *Nanophotonics* 8, 1327–1335, 2019

# Path-entanglement between topological modes



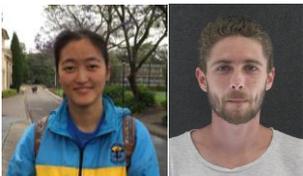
Michelle Wang Cooper Doyle

ZERO COUNTS  
in odd wgs

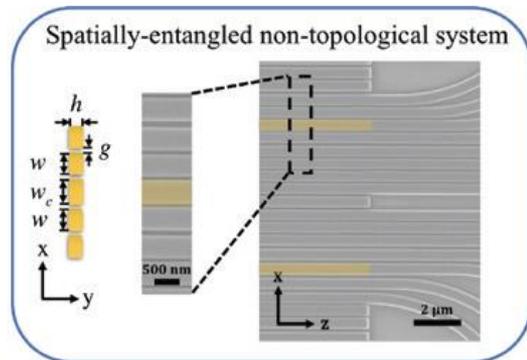


M. Wang, C. Doyle, B. Bell, M. J. Collins, E. Magi, B. J. Eggleton, M. Segev, and A. Blanco-Redondo, *Nanophotonics* 8, 1327–1335, 2019

# Path-entanglement between trivial modes

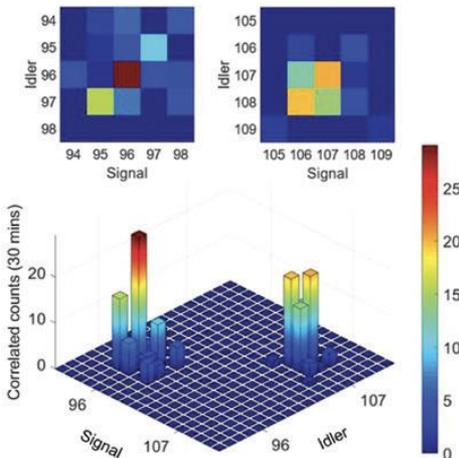


Michelle Wang Cooper Doyle



Trivial: Experiments

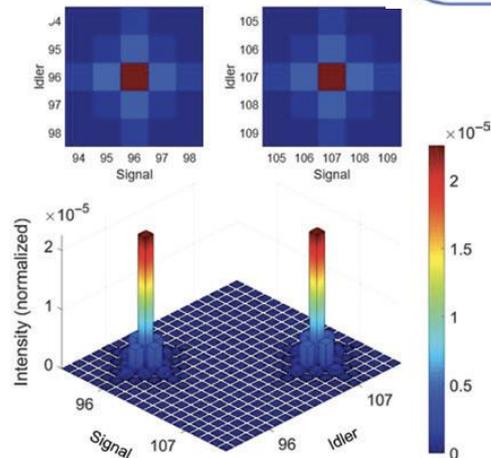
Disorder: 0%



No protected feature

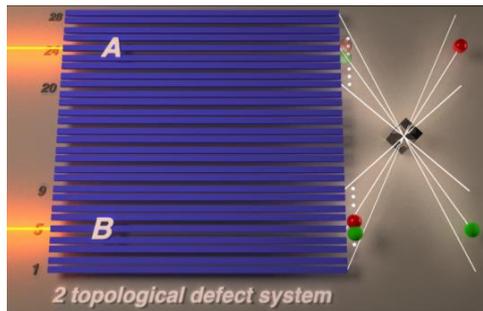
Trivial: Simulations

Disorder: 0%



M. Wang, C. Doyle, B. Bell, M. J. Collins, E. Magi, B. J. Eggleton, M. Segev, and A. Blanco-Redondo, *Nanophotonics* 8, 1327–1335, 2019

# Probing path-entanglement robustness by quantum interference (simulations)



Before the interferometer:

$$|\psi\rangle = \frac{|A\rangle_s |A\rangle_t + e^{2i\phi} |B\rangle_s |B\rangle_t}{\sqrt{2}}$$

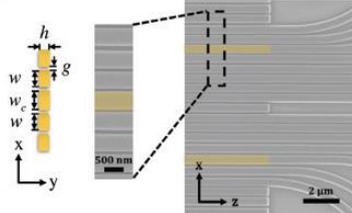
After the interferometer:

$$|\psi_{\text{bunch}}\rangle = \frac{i}{\sqrt{2}} (|B\rangle_s |B\rangle_t - |A\rangle_s |A\rangle_t)$$

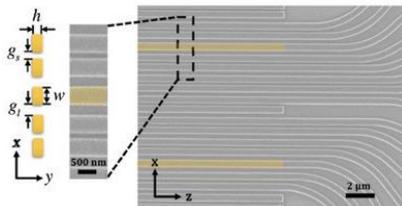
$$|\psi_{\text{split}}\rangle = \frac{i}{\sqrt{2}} (|A\rangle_s |B\rangle_t + |B\rangle_s |A\rangle_t)$$

$$|\psi_{\text{out}}\rangle = \cos\phi |\psi_{\text{split}}\rangle + \sin\phi |\psi_{\text{bunch}}\rangle$$

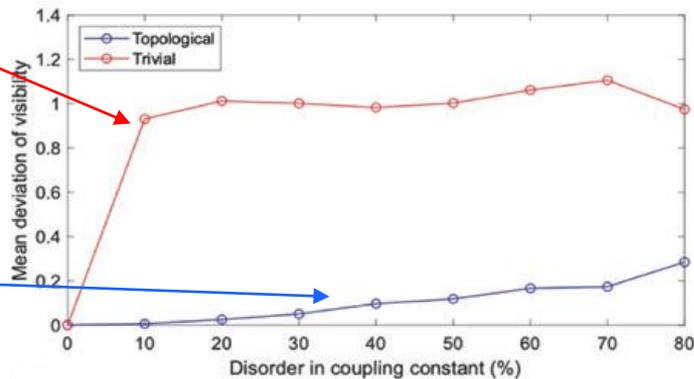
Spatially-entangled non-topological system



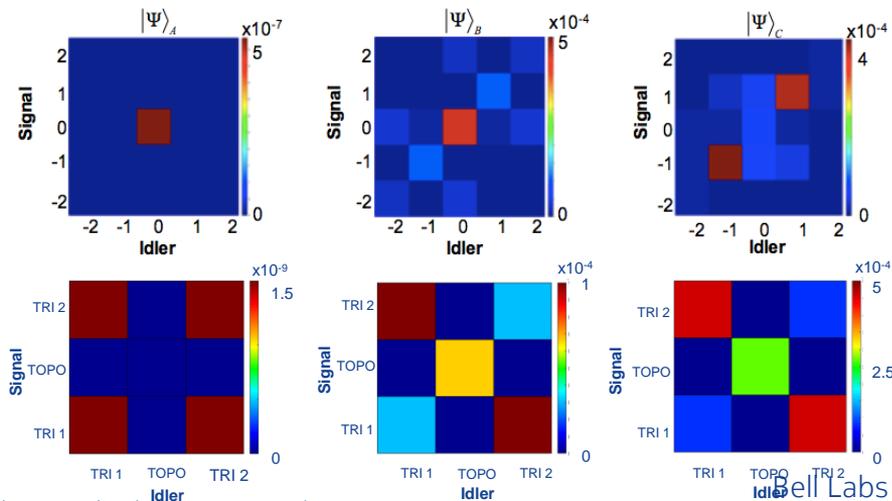
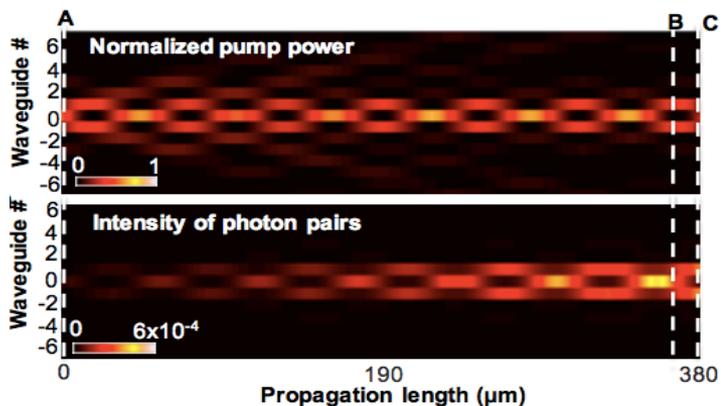
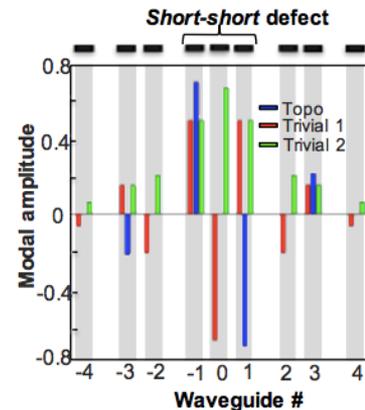
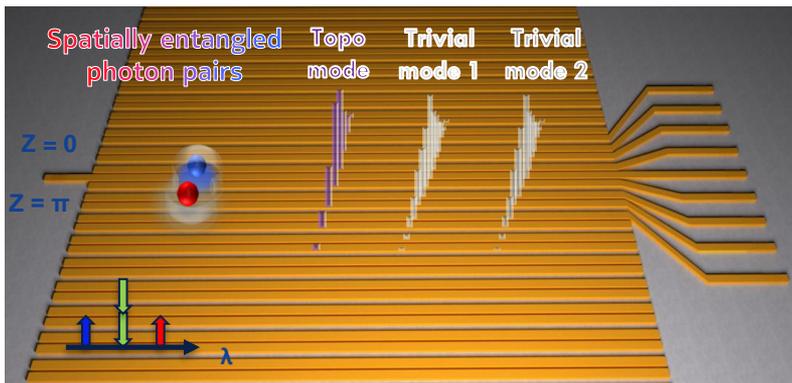
Spatially-entangled topological system



$$\text{Visibility} = \frac{\cos^2\phi - \sin^2\phi}{\cos^2\phi + \sin^2\phi} = \cos 2\phi$$

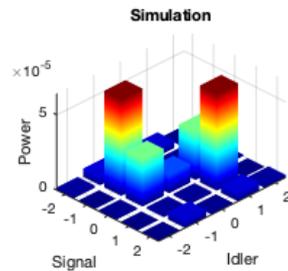
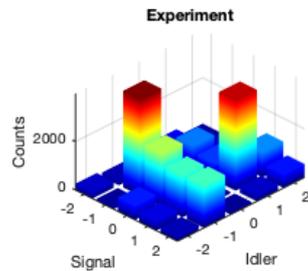


# Entanglement between topological and trivial modes

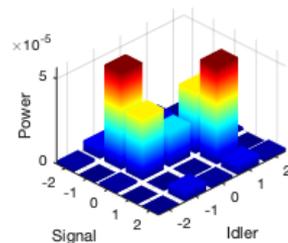
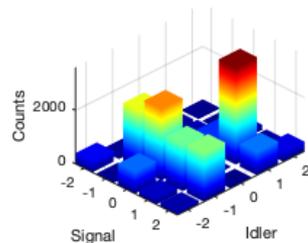


# Entanglement between topological and trivial modes

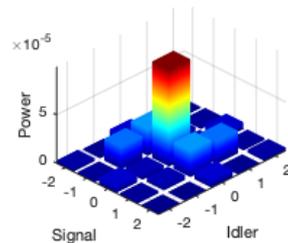
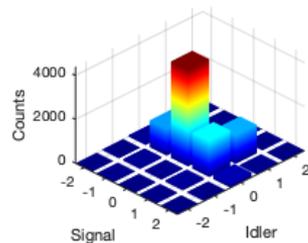
No disorder



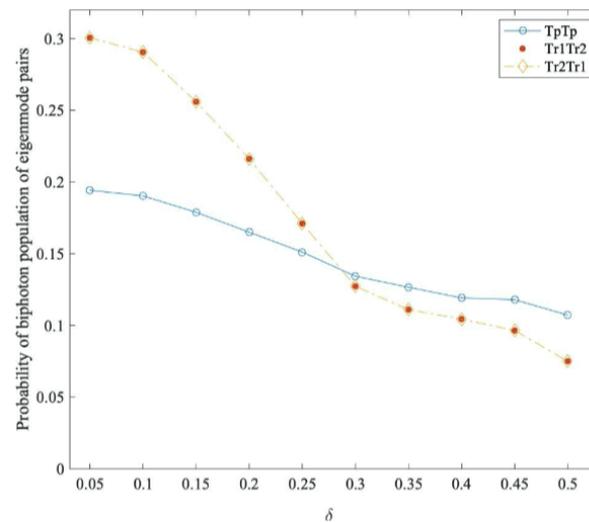
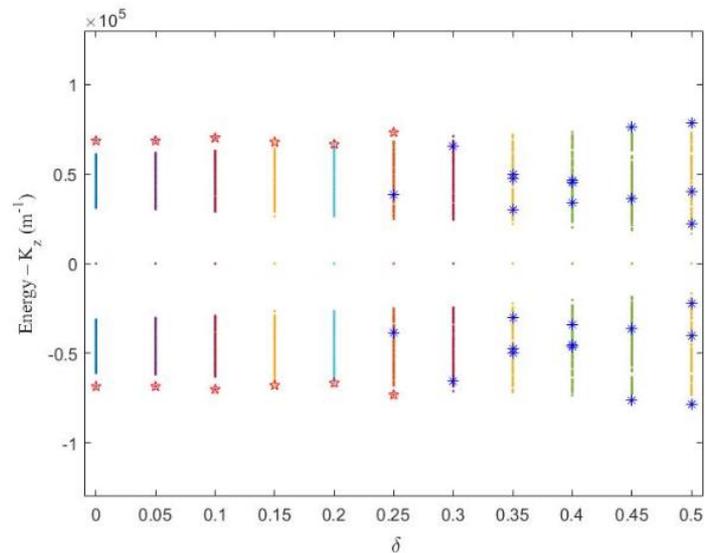
20% disorder



40% disorder



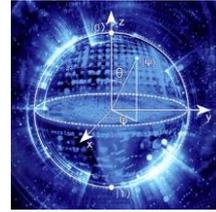
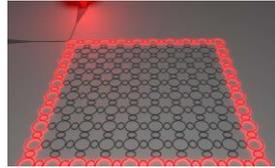
# Entanglement between topological and trivial modes



# End of Part 2: Conclusions

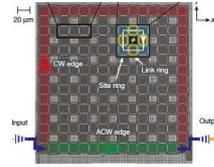
- ✓ Topological photonics is moving towards applications

Topological lasers    Topological quantum photonics

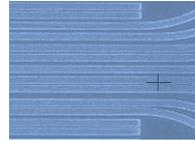


- ✓ Topology can protect quantum information

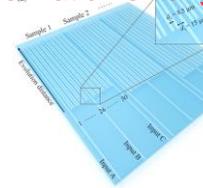
Spectral correlations



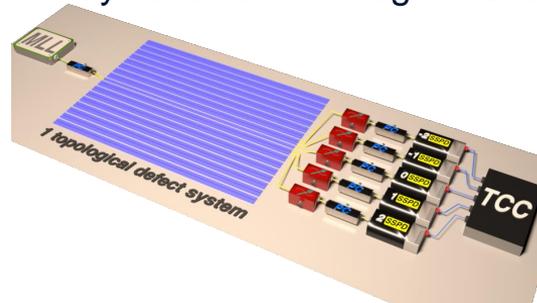
Path entanglement



Polarization entanglement



- ✓ Arrays of silicon waveguides to test topological quantum concepts



CMOS compatible

On-chip generation of quantum states

Convenient coupling, low loss, room temp. ...

# Thank you!

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